

Long time wave propagation in heterogeneous media

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CRC 1173

Wave
phenomena



Assyr Abdulle, Marcus J. Grote, and Christian Stohrer.

Finite element heterogeneous multiscale method for the wave equation: long-time effects.

Multiscale Model. Simul., 12(3):1230–1257, 2014.



Assyr Abdulle and Timothée Pouchon.

Effective models for the multidimensional wave equation in heterogeneous media over long time and numerical homogenization.

Math. Models Methods Appl. Sci., 26(14):2651–2684, 2016.



Assyr Abdulle and Timothée Pouchon.

A priori error analysis of the finite element heterogeneous multiscale method for the wave equation over long time.

SIAM J. Numer. Anal., 54(3):1507–1534, 2016.



Assyr Abdulle and Timothée Pouchon.

Effective models for long time wave propagation in locally periodic media.

SIAM J. Numer. Anal., 56(5):2701–2730, 2018.

Let $T^\varepsilon = \varepsilon^{-2}T$ and consider $u^\varepsilon : [0, T^\varepsilon] \times \Omega \rightarrow \mathbb{R}$ such that

$$\begin{aligned} \partial_t^2 u^\varepsilon(t, x) &= \partial_x (a^\varepsilon(x) \partial_x u^\varepsilon(t, x)) && \text{in } (0, T^\varepsilon] \times \Omega, \\ x \mapsto u^\varepsilon(t, x) &\Omega\text{-periodic} && \text{in } [0, T^\varepsilon], \\ u^\varepsilon(0, x) &= g^0(x), \quad \partial_t u^\varepsilon(0, x) = g^1(x) && \text{in } \Omega, \end{aligned}$$

where for $Y \subseteq \mathbb{R}$, $a^\varepsilon(x) = a\left(\frac{x}{\varepsilon}\right) = a(y)$ is Y -periodic in y and g^0, g^1 are given initial conditions.

$$\partial_t^2 u^0(t, x) = a^0 \partial_x^2 u^0(t, x)$$

where the effective coefficient $a^0 = \langle a(y)(1 + \partial\chi) \rangle_Y$ and χ is the periodic solution of an elliptic boundary value problem.

Effective model for time up to $\varepsilon^{-2}T$

$$\partial_t^2 \bar{u}(t, x) = a^0 \partial_x^2 \bar{u}(t, x) + \varepsilon^2 b^0 \partial_x^2 \partial_t^2 \bar{u}(t, x)$$

where the effective coefficient $a^0 = \langle a(y)(1 + \partial\chi) \rangle_Y$ and χ is the periodic solution of an elliptic boundary value problem.
 b^0 is obtained by a cascade of cell problems.

Agnes Lamacz, Dispersive effective models for waves in heterogeneous media, Math. Models Methods Appl. Sci., 21 (2011), pp. 1871–1899.

The solution \tilde{u} of

$$\partial_t^2 \tilde{u} = a^0 \partial_x^2 \tilde{u} - \varepsilon^2 \left(\tilde{a}^2 \partial_x^4 \tilde{u} - \tilde{b}^0 \partial_x^2 \partial_t^2 \tilde{u} \right)$$

approximates the true solution with error $\mathcal{O}(\varepsilon)$ up to time $\mathcal{O}(\varepsilon^{-2})$.

Tomas Dohnal, Agnes Lamacz, and Ben Schweizer, Bloch-wave homogenization on large time scales and dispersive effective wave equations, *Multiscale Model. Simul.*, 12 (2014), pp. 488–513.

Tomas Dohnal, Agnes Lamacz, and Ben Schweizer, Dispersive homogenized models and coefficient formulas for waves in general periodic media, *ArXiv e-prints*, (2014).

A family of effective equations for the wave equation over long time

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Equivalence class of correctors

Let $\chi \in \mathcal{W}_{\text{per}}(Y)$ be the unique (equivalence class of) solution of the cell problem

$$(a(y)\partial\chi, \partial\mathbf{w})_{L^2(Y)} = -(a(y), \partial\mathbf{w})_{L^2(Y)} \quad \text{for all } \mathbf{w} \in \mathcal{W}_{\text{per}}(Y)$$

$$a^0 = \langle a(y)(1 + \partial\chi) \rangle_Y$$

Let $\tilde{b}^0, \tilde{a}^2 \geq 0$ be non-negative. Find $\tilde{u} : [0, T^\varepsilon] \times \Omega \rightarrow \mathbb{R}$ such that

$$\partial_t^2 \tilde{u} = a^0 \partial_x^2 \tilde{u} - \varepsilon^2 \left(\tilde{a}^2 \partial_x^4 \tilde{u} - \tilde{b}^0 \partial_x^2 \partial_t^2 \tilde{u} \right) \quad \text{in } (0, T^\varepsilon] \times \Omega,$$

$$x \mapsto \tilde{u}(t, x) \text{ } \Omega\text{-periodic} \quad \text{in } [0, T^\varepsilon],$$

$$\tilde{u}(0, x) = g^0(x), \quad \partial_t \tilde{u}(0, x) = g^1(x) \quad \text{in } \Omega.$$

Theorem

Assume that the tensor $a^\varepsilon(x) = a\left(\frac{x}{\varepsilon}\right) = a(y)$ is uniformly Y -periodic and $a(y) \in W^{1,\infty}(Y)$. Furthermore, assume that the solution \tilde{u} of the homogeneous equation and the initial conditions satisfy the regularity

$$\tilde{u} \in L^\infty(0, T^\varepsilon; H^5(\Omega)), \quad \partial_t \tilde{u} \in L^\infty(0, T^\varepsilon; H^4(\Omega)), \quad \partial_t^2 \tilde{u} \in L^\infty(0, T^\varepsilon; H^3(\Omega))$$

$$g^0 \in H^4(\Omega), \quad g^1 \in H^3(\Omega).$$

Let χ be the solution of the cell problem and assume that for a $\chi \in \chi$, the coefficients \tilde{b}^0, \tilde{a}^2 satisfy the relation

$$a^0 \tilde{b}^0 - \tilde{a}^2 = a^0 \langle \chi^2 \rangle_Y - a^0 \langle \chi \rangle_Y^2$$

Then the following error estimate holds

$$\|u^\varepsilon - \tilde{u}\|_{L^\infty(0, T^\varepsilon; L^2(\Omega))} \leq C\varepsilon \left(\|g^1\|_{H^3(\Omega)} + \|g^0\|_{H^4(\Omega)} + \|\tilde{u}\|_{L^\infty(0, T^\varepsilon; H^5(\Omega))} + \|\partial_t^2 \tilde{u}\|_{L^\infty(0, T^\varepsilon; H^3(\Omega))} \right)$$

where C depends only on Ω, T, Y, a, λ and Λ .

$$\tilde{b}^0 = b^0 + \langle \chi \rangle_Y^2, \quad b^0 = \left\langle (\chi - \langle \chi \rangle_Y)^2 \right\rangle_Y, \quad \tilde{a}^2 = a^0 \langle \chi \rangle_Y^2$$

The set

$$\mathcal{F} = \left\{ \tilde{u}_{\langle \chi \rangle} \text{ solution of homogeneous equation with } \tilde{b}^0, \tilde{a}^2 \text{ defined above} \right\}$$

constitutes a family of effective solutions for u^ε .

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Now let $f \in L^2(0, T^\varepsilon; L^2_0(\Omega))$. Find $u^\varepsilon : [0, T^\varepsilon] \times \Omega \rightarrow \mathbb{R}$ such that

$$\begin{aligned} \partial_t^2 u^\varepsilon(t, x) - \partial_x (a^\varepsilon(x) \partial_x u^\varepsilon(t, x)) &= f(t, x) && \text{in } (0, T^\varepsilon] \times \Omega, \\ x \mapsto u^\varepsilon(t, x) &\Omega\text{-periodic} && \text{in } [0, T^\varepsilon], \\ u^\varepsilon(0, x) = g^0(x), \quad \partial_t u^\varepsilon(0, x) &= g^1(x) && \text{in } \Omega, \end{aligned}$$

Here we allow $a^\varepsilon(x) = a(x, \frac{x}{\varepsilon})$ where $a(x, y)$ is Ω -periodic in x and Y -periodic in y .

Among the family of effective equations \mathcal{F} we pick $\langle \chi \rangle_Y = 0$ which yields $\tilde{a}^2 = 0$ and denote the effective solution corresponding to this choice by \bar{u} .

Zero mean cell correctors and parameters

For each $x \in \Omega$ define $\chi(x, \cdot) \in W_{\text{per}}(Y)$ as the unique solution of

$$(a(x, \cdot) \partial_y \chi(x, \cdot), \partial_y w)_{L^2(Y)} = -(a(x, \cdot), \partial_y w)_{L^2(Y)} \quad \text{for all } w \in W_{\text{per}}(Y)$$

For $x \in \Omega$ define

$$a^0(x) = \langle a(x, \cdot) (1 + \partial_y \chi(x, \cdot)) \rangle_Y, \quad b^0(x) = \langle (\chi(x, \cdot))^2 \rangle_Y$$

The equation for \bar{u} is: Find $\bar{u} : [0, T^\varepsilon] \times \Omega \rightarrow \mathbb{R}$ such that

$$\begin{aligned} \partial_t^2 \bar{u} - \partial_x \left(a^0(x) \partial_x \bar{u} \right) - \varepsilon^2 \partial_x \left(b^0(x) \partial_x \partial_t^2 \bar{u} \right) &= f && \text{in } (0, T^\varepsilon] \times \Omega, \\ x \mapsto \bar{u}(t, x) \text{ } \Omega\text{-periodic} &&& \text{in } [0, T^\varepsilon], \\ \bar{u}(0, x) = g^0(x), \quad \partial_t \bar{u}(0, x) = g^1(x) &&& \text{in } \Omega. \end{aligned}$$

Define bilinear forms

$$A^0(v, w) = \left(a^0(x) \partial_x v, \partial_x w \right)_{L^2(\Omega)}, \quad B^0(v, w) = \left(b^0(x) \partial_x v, \partial_x w \right)_{L^2(\Omega)}$$

and the Hilbert space $\mathcal{S}(\Omega) = \{v \in L^2(\Omega) : \partial_x v \in L^2(\Omega)\}$ equipped with

$$(v, w)_{\mathcal{S}(\Omega)} = (v, w)_{L^2(\Omega)} + \varepsilon^2 B^0(v, w)$$

\bar{u} satisfies

$$\left(\partial_t^2 \bar{u}(t), v \right)_{\mathcal{S}(\Omega)} + A^0(\bar{u}(t), v) = (f(t), v)_{L^2(\Omega)}, \quad \forall v \in W_{\text{per}}(\Omega), \text{ a.e. } t$$

$$\bar{u}(0) = g^0, \quad \partial_t \bar{u}(0) = g^1$$

Find $u_H : [0, T^\varepsilon] \rightarrow V_H(\Omega)$ such that

$$\left(\partial_t^2 u_H(t), v_H \right)_Q + A_H(u_H(t), v_H) = (f(t), v_H)_{L^2(\Omega)}, \quad \forall v_H \in V_H(\Omega), \text{ a.e. } t$$

$$u_H(0) = g_H^0, \quad \partial_t u_H(0) = g_H^1$$

where

$$A_H(v_H, w_H) = \sum_{K \in \mathcal{T}_H} \sum_{j=1}^J \omega_{K_j} a_K^0(x_{K_j}) \partial_x v_H(x_{K_j}) \partial_x w_H(x_{K_j})$$

$$(v_H, w_H)_Q = (v_H, w_H)_H + (v_H, w_H)_M,$$

$$(v_H, w_H)_H = \sum_{K \in \mathcal{T}_H} \sum_{j=1}^{J'} \omega'_{K_j} v_H(x'_{K_j}) w_H(x'_{K_j}),$$

$$(v_H, w_H)_M = \varepsilon^2 \sum_{K \in \mathcal{T}_H} \sum_{j=1}^J \omega_{K_j} b_K^0(x_{K_j}) \partial_x v_H(x_{K_j}) \partial_x w_H(x_{K_j})$$

Let $\psi_{h,K_j} \in V_h(K_{\delta_j})$ the solution of

$$\left(a^\varepsilon(x) \partial_x \psi_{h,K_j}, \partial_x z_h \right)_{L^2(K_{\delta_j})} = - \left(a^\varepsilon(x), \partial_x z_h \right)_{L^2(K_{\delta_j})}, \quad \forall z_h \in V_h(K_{\delta_j})$$

and define approximated parameters at quadrature points as

$$a_K^0(x_{K_j}) = \left\langle a^\varepsilon(x) (1 + \partial_x \psi_{h,K_j}) \right\rangle_{K_{\delta_j}}, \quad b_K^0(x_{K_j}) = \varepsilon^{-2} \left\langle \left(\psi_{h,K_j} \right)^2 \right\rangle_{K_{\delta_j}}$$

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A priori $L^\infty(H^1)$ error

Theorem

Assume that δ satisfies $\frac{\delta}{\varepsilon} \in \mathbb{N}_{>0}$, that the micro mesh size is $h < \varepsilon$ and that the degree of the micro finite element space is $q = 1$. Furthermore, assume that the tensor $a^\varepsilon \in W^{1,\infty}(\Omega)$ is collocated in the slow variable, i.e. for all $(K, j) \in \mathcal{T}_H \times \{1, \dots, J\}$, $a^\varepsilon(x) = a\left(x_{K_j}, \frac{x}{\varepsilon}\right)$ for a.e. $x \in K_{\delta_j}$. Finally, assume that $a^0, b^0 \in W^{\ell,\infty}(\Omega)$ and $\partial_t^k \bar{u} \in L^\infty(0, T^\varepsilon; H^{\ell+1}(\Omega))$ for $0 \leq k \leq 4$. Then the error $e = \bar{u} - u_H$ satisfies the estimate

$$\|\partial_t e\|_{L^\infty(0, T^\varepsilon; L^2(\Omega))} + \|e\|_{L^\infty(0, T^\varepsilon; H^1(\Omega))} \leq C_1 \left(h/\varepsilon^2\right)^2 + e_{H^1}^{FE},$$

where $e_{H^1}^{FE}$ is the standard FEM error estimate,

$$e_{H^1}^{FE} \leq C_2 \left(\|g^1 - g_H^1\|_{L^2(\Omega)} + \varepsilon \|g^1 - g_H^1\|_{H^1(\Omega)} + \|g^0 - g_H^0\|_{H^1(\Omega)} + H^\ell \right),$$

$C_1 = \tilde{C}_1 \sum_{k=0}^4 \|\partial_t^k \bar{u}\|_{L^\infty(H^{\ell+1}(\Omega))}$ and $C_2 = \tilde{C}_2 \sum_{k=0}^4 \|\partial_t^k \bar{u}\|_{L^1(H^{\ell+1}(\Omega))}$ with \tilde{C}_1, \tilde{C}_2 independent of H, h, ε and δ .

A priori $L^\infty(L^2)$ error

Theorem

As before assume that $h \leq \varepsilon$, $q = 1$ and $a^\varepsilon(x) = a\left(x_{K_j}, \frac{x}{\varepsilon}\right)$ for a.e. $x \in K_{\delta_j}$. Furthermore, assume that $\partial_t^k \bar{u} \in L^\infty(0, T^\varepsilon; H^{\ell+1}(\Omega))$ for $0 \leq k \leq 3$, $g_H^1 = I_H g^1$ and $a^0 \in W^{\ell+1, \infty}(\Omega)$. Then $e = \bar{u} - u_H$ satisfies the estimate

$$\|e\|_{L^\infty(0, T^\varepsilon; L^2(\Omega))} \leq C_1 \left(h/\varepsilon^2\right)^2 + e_{L^2}^{FE},$$

where $e_{L^2}^{FE}$ is the standard FEM error estimate,

$$e_{L^2}^{FE} \leq C_2 \left(\|g^0 - g_H^0\|_{L^2(\Omega)} + \varepsilon \|g^0 - g_H^0\|_{H^1(\Omega)} + H^{\ell+1} + \varepsilon H^\ell \right),$$

$C_1 = \tilde{C}_1 \sum_{k=0}^3 \|\partial_t^k \bar{u}\|_{L^\infty(H^{\ell+1}(\Omega))}$ and $C_2 = \tilde{C}_2 \sum_{k=0}^3 \|\partial_t^k \bar{u}\|_{L^1(H^{\ell+1}(\Omega))}$ with \tilde{C}_1, \tilde{C}_2 independent of H, h, ε and δ .

Corollary

Assume that the tensor $a^\varepsilon(x) = a\left(\frac{x}{\varepsilon}\right)$ is periodic with $a \in W^{1,\infty}(\Omega)$ and $f = 0$. Also, assume that $g_H^0 = I_H g^0$, $g_H^1 = I_H g^1$ and let the settings of the FE-HMM-L be such that $\delta/\varepsilon \in \mathbb{N}_{>0}$, $h \leq \varepsilon$, $q = 1$ and $\ell = 1$. Finally assume that the following regularity holds

$$g^0 \in H^4(\Omega), \quad g^1 \in H^3(\Omega), \quad \partial_t^k \in L^\infty(H^{5-k}(\Omega)) \quad 0 \leq k \leq 3.$$

Then we have the following estimate

$$\|u^\varepsilon - u_H\|_{L^\infty(0,T^\varepsilon;L^2(\Omega))} \leq C_1 \left(\varepsilon + \left(h/\varepsilon^2 \right)^2 \right) + C_2 \left(H^2 + \varepsilon H \right)$$

where $C_1 = \tilde{C}_1 \sum_{k=0}^3 \|\partial_t^k \tilde{u}\|_{L^\infty(H^{5-k}(\Omega))}$ and $C_2 = \tilde{C}_2 \sum_{k=0}^3 \|\partial_t^k \tilde{u}\|_{L^1(H^2(\Omega))}$ independent of H, h, ε and δ .

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$$\mathcal{B}^\varepsilon \tilde{u}(t, x) = \tilde{u}(t, x) + \varepsilon \chi \left(\frac{x}{\varepsilon} \right) \partial_x \tilde{u}(t, x) + \varepsilon^2 \theta \left(\frac{x}{\varepsilon} \right) \partial_x^2 \tilde{u}(t, x) + \varepsilon^3 \kappa \left(\frac{x}{\varepsilon} \right) \partial_x^3 \tilde{u}(t, x) + \varepsilon^4 \rho \left(\frac{x}{\varepsilon} \right) \partial_x^4 \tilde{u}(t, x)$$

$$\mathcal{A}^\varepsilon = -\partial_x \left(a \left(\frac{x}{\varepsilon} \right) \partial_x (\cdot) \right)$$

Now we compute $(\partial_t^2 + \mathcal{A}^\varepsilon) \mathcal{B}^\varepsilon \tilde{u}$

$$\partial_t^2 (\mathcal{B}^\varepsilon \tilde{u}) = a^0 \partial_x^2 \tilde{u} + \varepsilon a^0 \chi \partial_x^3 \tilde{u} + \varepsilon^2 \left(a^0 (\theta + \tilde{b}^0) - \tilde{a}^2 \right) \partial_x^4 \tilde{u} + \mathcal{O}(\varepsilon^3)$$

$$-\partial_x \left(a \left(\frac{x}{\varepsilon} \right) \partial_x (\mathcal{B}^\varepsilon \tilde{u}(x)) \right) = \varepsilon^{-1} \left(-\partial_y (a(y)(1 + \partial_y \chi)) \right) \partial_x \tilde{u}$$

$$\varepsilon^0 \left(-\partial_y (a(y)(\chi + \partial_y \theta)) - a(y)(1 + \partial_y \chi) \right) \partial_x^2 \tilde{u}$$

$$\varepsilon^1 \left(-\partial_y (a(y)(\theta + \partial_y \kappa)) - a(y)(\chi + \partial_y \theta) \right) \partial_x^3 \tilde{u}$$

$$\varepsilon^2 \left(-\partial_y (a(y)(\kappa + \partial_y \rho)) - a(y)(\theta + \partial_y \kappa) \right) \partial_x^4 \tilde{u} + \mathcal{O}(\varepsilon^3)$$

Lemma

Under the assumptions of Theorem, $\hat{\mathcal{B}}^\varepsilon \tilde{u}$ satisfies

$$\left(\partial_t^2 + \mathcal{A}^\varepsilon\right) \hat{\mathcal{B}}^\varepsilon \tilde{u}(t) = \mathcal{R}^\varepsilon \tilde{u}(t) \quad \text{in } \mathcal{W}_{\text{per}}^*(\Omega) \quad \text{for a.e. } t \in [0, T^\varepsilon],$$

where the right hand side $\mathcal{R}^\varepsilon \tilde{u} \in L^2(0, T^\varepsilon; \mathcal{W}_{\text{per}}^*(\Omega))$ satisfies the estimate

$$\|\mathcal{R}^\varepsilon \tilde{u}\|_{L^2(0, T^\varepsilon; \mathcal{W}_{\text{per}}^*(\Omega))} \leq C\varepsilon^2 \left(\|\tilde{u}\|_{L^\infty(0, T^\varepsilon; H^5(\Omega))} + \|\partial_t^2 \tilde{u}\|_{L^\infty(0, T^\varepsilon; H^3(\Omega))} \right)$$

for a constant C that only depends on T, Y, a, λ and Λ .

Lemma

Under the assumptions of Theorem, $\eta^\varepsilon = \hat{\mathcal{B}}^\varepsilon \tilde{u} - [u^\varepsilon]$ satisfies

$$\begin{aligned} & \|\partial_t \eta^\varepsilon\|_{L^\infty(\mathcal{W}_{\text{per}}^*(\Omega))} + \|\eta^\varepsilon\|_{L^\infty(\mathcal{L}^2)} \\ & \leq C \left(\|\partial_t \eta^\varepsilon(0)\|_{\mathcal{W}_{\text{per}}^*(\Omega)} + \|\eta^\varepsilon(0)\|_{\mathcal{L}^2} + \varepsilon^{-1} \|\mathcal{R}^\varepsilon \tilde{u}\|_{L^2(\mathcal{W}_{\text{per}}^*(\Omega))} \right), \end{aligned}$$

where C depends only on λ, Λ and T and $\mathcal{R}^\varepsilon \tilde{u}$ is given above.

Lemma

$$e_{a^0} = \sup_{K,j} |a^0(x_{K_j}) - a_K^0(x_{K_j})| \leq C \left(\frac{h}{\varepsilon}\right)^2, \quad e_{b^0} = \sup_{K,j} \varepsilon^2 |a^0(x_{K_j}) - a_K^0(x_{K_j})| \leq C\varepsilon \left(\frac{h}{\varepsilon}\right)^2$$